Some solitary wave solutions of generalized Pochhammer-Chree equation via Exp-function method

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Abstract—In this paper, Exp-function method is used for some exact solitary solutions of the generalized Pochhammer-Chree equation. It has been shown that the Exp-function method, with the help of symbolic computation, provides a very effective and powerful mathematical tool for solving nonlinear partial differential equations. As a result, some exact solitary solutions are obtained. It is shown that the Exp-function method is direct, effective, succinct and can be used for many other nonlinear partial differential equations.

Keywords—Exp-function method, Generalized Pochhammer-Chree equation, solitary wave solution, ODE’s

I. INTRODUCTION

The study of exact solutions of nonlinear partial differential equations (NPDE) plays an important role in mathematical physics, engineering and the other sciences. In the past several decades, various methods for obtaining solutions of NPDE’s and ODE’s have been presented, such as, tanh-function method [1], [2], [3], Adomian decomposition method [4], [5], Homotopy perturbation method [6], [7], [8], variational iteration method [9], [10], [11], spectral method [12], [13], [14], sine-cosine method [15], [16], radial basis method [17], [18] and so on. Recently, Ji-Huan He and Xia-Hong Wu [19] proposed a novel method, so called Exp-function method, which is easy, succinct and powerful to implement to nonlinear partial differential equations arising in mathematical physics. The Exp-function method has been successfully applied to many kinds of NPDEs, such as, KdV equation with variable coefficients [20], Kawahara equation [22], Boussinesq equations [23], Burger’s equations [24], [25], [26], Double Sine-Gordon equation [27], [28], Fisher equation [29], Jaulent-Miodek equations [30] and the other important nonlinear partial differential equations [31], [32], [33]. In this paper we apply the Exp-function method [19] to obtain exact solitary wave solution of a nonlinear partial differential equation, namely, generalized Pochhammer-Chree equation (GPC) given by

\[ u_{tt} - u_{txxx} - (\alpha u + \beta u^{n+1} + \gamma u^{2n+1})_{xx} = 0, \quad n \geq 1. \]

where \( \alpha, \beta \) and \( \gamma \) are constants. GPC equation represents a nonlinear model of longitudinal wave propagation of elastic rods [34], [35], [36], [37], [38], [39], [40], [41], [42], [43], [44], [45], [46]. The model for \( \alpha = 1, \beta = \frac{1}{n+1} \) and \( \gamma = 0 \)

was studied in [40], [41] where solitary wave solutions for this model was obtained for \( n = 1, 2 \) and 4. A second model for \( \alpha = 0, \beta = -\frac{1}{2} \) and \( \gamma = 0 \) was studied by [42] and solitary wave solutions were obtained as well. However, a third model was investigated in [37], [43], [44], [45], [46] for \( n = 1 \) and \( n = 2 \) where explicit solitary wave solutions and kinks solutions were derived.

The rest of the paper is organized as follows: Section 2 describes exp-function method for finding exact solutions to the NPDEs. The applications of the proposed analytical scheme presented in Section 3. The conclusions are discussed in the section 4. Exp-function calculations are provided in the end.

II. BASIC IDEA OF EXP-FUNCTION METHOD

We consider a general nonlinear PDE in the following form

\[ N(u, u_x, u_t, u_{xx}, u_{tt}, u_{xt}, \ldots) = 0, \quad (1) \]

where \( N \) is a polynomial function with respect to the indicated variables or some functions which can be reduced to a polynomial function by using some transformation. We introduce a complex variation as

\[ u(x,t) = U(\eta), \quad \eta = k(x - ct) + \varphi_0. \quad (2) \]

where \( k \) and \( c \) are constants and \( \varphi_0 \) is an arbitrary constant. We can rewrite Eq.(1) in the following nonlinear ordinary differential equations

\[ N(U, ku', -kcU', k^2u'', \ldots) = 0, \]

where the prime denotes the derivation with respect to \( \eta \). According to the Exp-function method [19], we assume that the solution can be expressed in the form

\[ U(\eta) = \sum_{i=-d}^{c} a_i \exp(i \eta) + \sum_{j=-q}^{p} b_j \exp(j \eta), \quad (3) \]

where \( c, d, p, q \) and \( q \) are positive integers which can be freely chosen, \( a_i \) and \( b_j \) are unknown constants to be determined. To determine the values of \( d \) and \( q \), we balance the highest order linear term with the highest order nonlinear term in Eq.(3). Similarly to determine the values of \( d \) and \( q \). So by means of the exp-function method, we obtain the generalized solitary solution and periodic solution for nonlinear evolution equations arising in mathematical physics.
III. APPLICATIONS OF THE EXP-FUNCTION METHOD

In this section, we show the detailed steps of the Exp-function method to construct exact solitary wave solutions of generalized Pochhammer-Chree equations (GPC)

\[ u_{tt} - u_{ttxx} - (\alpha u + \beta u^{n+1} + \gamma u^{2n+1})_{xx} = 0 , \]  

where \( \alpha, \beta \) and \( \gamma \) are constants. Making the travelling wave transformation

\[ u(x, t) = U(\eta), \quad \eta = k(x - ct) + \varphi_0 , \]

and integrating twice, here \( k \) and \( c \) are constants to be determined later, then Eq.(4) becomes an ordinary differential equation in the form

\[ k^2(c^2 - \alpha)U - k^4c^2U'' - k^2\beta U^{n+1} - k^2\gamma U^{2n+1} = 0 , \]

where the prime denotes the derivative with respect to \( \eta \) and also where the integration constants are chosen as zero. We now use the transformation

\[ U^n = v , \]

which we find

\[ U'' = \frac{1 - n}{n^2} v^{\frac{1}{n} - 2} (v')^2 + \frac{1}{n} v^{\frac{1}{n} - 1} v'' , \]

substituting the transformations (5) into the GPC equation gives the ODE,

\[ n^2k^2(c^2 - \alpha)v^2 - k^4c^2(1 - n)(v')^2 - k^4c^2nvv'' - k^2\beta n^2 v^3 - k^2\gamma n^2 v^4 = 0 , \]  

which leads to results

\[ -(3q + 2d) = -(q + 4d) , \]

which reduces to

\[ q = d . \]

For simplicity, we set \( p = c = 1 \) and \( q = d = 1 \), so Eq.(3) reduces to

\[ v(\eta) = a_1 \exp(\eta) + a_0 + a_{-1} \exp(-\eta) \ exp(\eta) + b_0 + b_{-1} \exp(-\eta) . \]

Similarly to determine values of \( d \) and \( q \), we balance the linear term of lowest order in Eq.(6)

\[ vv'' = \frac{... + d_1 \exp(-(3q + 2d)\eta)}{... + d_2 \exp(-5q\eta)} , \]  

and

\[ v^4 = \frac{... + d_1 \exp(-(q + 4d)\eta)}{... + d_2 \exp(-5q\eta)} , \]

where \( d_i \) are determined coefficients only for simplicity, we have

\[ (3q + 2d) = -(q + 4d) , \]

Substituting Eq.(11) into Eq.(6), equating to zero the coefficients of all powers of \( \exp(n\eta) \) yields a set of algebraic equations for \( a_0, b_0, a_{-1}, a_1, b_1, k \) and \( c \) (see Appendix A). By solving the system of algebraic equations with a professional mathematical software, we obtain

\[ a_{1} = 0 , \quad a_{0} = \frac{b_{0}}{\beta} (c^2 - \alpha)(n + 2) , \quad a_{-1} = 0 , \]

\[ b_{0} = b_{0} , \quad k = \frac{b_{0}}{c} \sqrt{c^2 - \alpha} , \quad c = c , \]

\[ b_{-1} = \frac{b_{1}^{2}}{4\beta^2(n + 1)} \left[ \gamma(c^2 - \alpha)(n + 2)^2 + \beta^2(n + 1) \right] . \]

Substituting these result into Eq.(11), we obtain

\[ v(\eta) = \frac{b_{0}}{\beta} (c^2 - \alpha)(n + 2) \ exp(\eta) + b_{0} + b_{-1} \exp(-\eta) . \]

where \( b_0 \) and \( c \) are free parameters and

\[ b_{-1} = \frac{b_{0}^{2}}{4\beta^2(n + 1)} \left[ \gamma(c^2 - \alpha)(n + 2)^2 + \beta^2(n + 1) \right] . \]

To compare our results with those obtained in [43], [45], if we set

\[ b_{0} = \frac{2\sqrt{3}\beta}{\sqrt{3\beta^2 + 16\gamma(c^2 - \alpha)}} , \quad n = 2 , \]

Eq(12) becomes

\[ v(\eta) = \frac{8\sqrt{\frac{3}{\sqrt{3\beta^2 + 16\gamma(c^2 - \alpha)}}} (c^2 - \alpha)}{\exp(\eta) + \frac{2\sqrt{3}\beta}{\sqrt{3\beta^2 + 16\gamma(c^2 - \alpha)}} \exp(-\eta)} . \]
where \( \eta = \frac{2}{c} \sqrt{c^2 - \alpha} (x - ct) + \varphi_0 \). We re-write Eq.(13) and use of Eq.(5) in the form

\[
\begin{align*}
\eta &= \frac{2}{c} \sqrt{c^2 - \alpha} (x - ct) + \varphi_0 \\
\text{which is the traveling wave solution obtained in [43], [45].} \\
\text{Also, by the choice } \gamma = 0 \text{ in our solution (12) gives}
\end{align*}
\]

\[
v(\eta) = \frac{2(2c^2 - \alpha)(n + 2)}{\exp(\eta) + b_0 + \frac{\beta}{\gamma^2} \exp(-\eta)} .
\]

(IV) At \( c^2 < \alpha \) and \( b_0 = -2 \).

We can obtain from Eq.(14) and Eq.(5) that

\[
\eta = \frac{2}{c} \sqrt{c^2 - \alpha} (x - ct) + \varphi_0 \\
\text{or equivalently}
\]

\[
\begin{align*}
u(\eta) &= \frac{2(2c^2 - \alpha)(n + 2)}{\exp(\eta) + b_0 + \frac{\beta}{\gamma^2} \exp(-\eta)} .
\end{align*}
\]

(II) At \( c^2 > \alpha \) and \( b_0 = -2 \).

We can obtain from Eq.(14) and Eq.(5) that

\[
\eta = \frac{2}{c} \sqrt{c^2 - \alpha} (x - ct) + \varphi_0
\]

or equivalently

\[
\eta = \frac{2}{c} \sqrt{c^2 - \alpha} (x - ct) + \varphi_0
\]

(III) At \( c^2 < \alpha \) and \( b_0 = 2 \).

We can obtain from Eq.(14) and Eq.(5) that

\[
\eta = \frac{2}{c} \sqrt{c^2 - \alpha} (x - ct) + \varphi_0
\]

or equivalently
We can obtain from Eq. (17) and Eq. (5) that

\[
u(x,t) = \left\{ \begin{array}{l}
\frac{1}{2} \left[ \frac{(n+1)(\alpha - c^2)}{\gamma} \right] \\
\tanh \left[ \frac{n}{2c} \sqrt{c^2 - \alpha} (x - ct) + \varphi_0 \right] \\
- \coth \left[ \frac{n}{2c} \sqrt{c^2 - \alpha} (x - ct) + \varphi_0 \right]
\end{array} \right\}^{1/n}
\]


\[ 4n^2 c_{2} a_{1} a_{0} b_{0} - 4n^2 c_{2} a_{1} a_{0} b_{0} + 4k^2 c_{2} n_{1} a_{1} a_{0} - 4k^2 c_{2} n_{1} a_{1} a_{0} b_{1} - 4n^2 c_{2} n_{1} a_{1} a_{0} b_{1} - 4n^2 c_{2} n_{1} a_{1} a_{0} b_{1} = 0 , \]

\[ - n^2 c_{2} a_{1} b_{0} + n^2 c_{2} a_{1} b_{0} - n^2 c_{2} a_{1} b_{0} - n^2 c_{2} a_{1} b_{0} - n^2 c_{2} a_{1} b_{0} - n^2 c_{2} a_{1} b_{0} = 0 , \]

\[ + 2n^2 c_{2} a_{1} b_{0} + 2n^2 c_{2} a_{1} b_{0} + 2n^2 c_{2} a_{1} b_{0} + 2n^2 c_{2} a_{1} b_{0} + 2n^2 c_{2} a_{1} b_{0} = 0 , \]

\[ + n^2 c_{2} a_{1} b_{0} + n^2 c_{2} a_{1} b_{0} + 2n^2 c_{2} a_{1} b_{0} + 2n^2 c_{2} a_{1} b_{0} + 2n^2 c_{2} a_{1} b_{0} = 0 , \]

\[ - n^2 c_{2} a_{1} b_{0} - n^2 c_{2} a_{1} b_{0} - n^2 c_{2} a_{1} b_{0} - n^2 c_{2} a_{1} b_{0} - n^2 c_{2} a_{1} b_{0} = 0 , \]

\[ - 2n^2 c_{2} a_{1} b_{0} - 2n^2 c_{2} a_{1} b_{0} - 2n^2 c_{2} a_{1} b_{0} - 2n^2 c_{2} a_{1} b_{0} - 2n^2 c_{2} a_{1} b_{0} = 0 , \]

\[ - n^2 c_{2} a_{1} b_{0} - n^2 c_{2} a_{1} b_{0} - n^2 c_{2} a_{1} b_{0} - n^2 c_{2} a_{1} b_{0} - n^2 c_{2} a_{1} b_{0} = 0 , \]

\[ - k^2 c_{2} n_{1} a_{0} b_{0} + nk^2 c_{2} a_{0} - 4n^2 c_{2} a_{0} b_{0} - 2n^2 c_{2} a_{0} b_{0} + 2n^2 c_{2} a_{0} b_{0} + 2n^2 c_{2} a_{0} b_{0} + 2n^2 c_{2} a_{0} b_{0} + 2n^2 c_{2} a_{0} b_{0} = 0 , \]

\[ - k^2 c_{2} n_{1} a_{0} b_{0} + nk^2 c_{2} a_{0} - 4n^2 c_{2} a_{0} b_{0} - 2n^2 c_{2} a_{0} b_{0} + 2n^2 c_{2} a_{0} b_{0} + 2n^2 c_{2} a_{0} b_{0} + 2n^2 c_{2} a_{0} b_{0} + 2n^2 c_{2} a_{0} b_{0} = 0 , \]

\[ - n^2 c_{2} a_{0} b_{0} - n^2 c_{2} a_{0} b_{0} - n^2 c_{2} a_{0} b_{0} - n^2 c_{2} a_{0} b_{0} - n^2 c_{2} a_{0} b_{0} = 0 , \]

\[ - n^2 c_{2} a_{0} b_{0} - n^2 c_{2} a_{0} b_{0} - n^2 c_{2} a_{0} b_{0} - n^2 c_{2} a_{0} b_{0} - n^2 c_{2} a_{0} b_{0} = 0 , \]

\[ \text{REFERENCES} \]


